



ANALYTICAL AND NUMERICAL INVESTIGATION OF FREE VIBRATION OF NANOPARTICLE-REINFORCED COMPOSITE CYLINDRICAL SHELLS

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Abstract

This study focuses on the characterization of free vibration of composite shell structures strengthened by different volume fractions of nanoparticles analytically and numerically. Using simply supported boundary conditions, the governing differential equation of motion for the shell was formulated based on the Donnell-Mushtari-Vlasov (DMV) shell theory. For different design parameters, the natural frequency was investigated by employing the Orthogonality method. Four different layers of material, namely Perlon, Carbon, Kevlar, and Kenaf, of thickness 30 mm, were made. Nanoparticles Alumina (Al_2O_3) and Silica (SiO_2) were chosen and mixed in varying volume fractions (0.5%, 1%, 1.5%, 2%, and 2.5%) for the sample fabrication. Two types of samples, A and B, were created based on the arrangement of layers. The tensile tests were performed on the fabricated specimens to identify the longitudinal Young's modulus of specimens. The two groups that consist of different layers of materials were made and named as group A and group B. The results indicate an increase in Young's modulus of 33.9% increase for nano Al_2O_3 and a 42.25% increase for nano SiO_2 at a volume fraction of 2.5% for group A, while for group B, the enhancement was 37.96% and 47.39% for Al_2O_3 and SiO_2 nanoparticles, respectively. The results indicate that as the volume fraction of nanomaterial is increased, the natural frequency increases. The experimental results are used to validate both analytical and numerical solution conducted by the finite element method (FEM) under various loading conditions. The maximum difference between the analytical and numerical prediction of the natural frequency results was within 5%.

Keywords: cylindrical shell structure, free vibration frequency, nanoparticles, FEM

NOMENCLATURE

R	The radius of the cylinder (mm)
L	The length of the shell (mm)
V_f	Volume fraction %
E	Modulus of elasticity (GPa)
ν	Poisson's ratio
ρ	The mass density of shell material (kg/m^3)
ω	Natural frequency (rad/sec)
ψ	Non-dimensional parameter of the natural frequency

1. INTRODUCTION

Compared with classical designs, nanocomposite structures exhibit unusual property

combinations and can be designed in various ways with high-performance materials [1-3]. Curved surfaces characterize structural membranes or shells and can distribute loads in multiple directions. Properly constructed, shaped, proportioned, and supported, they maintain structural integrity without bending or twisting. Thin shells specifically refer to shells with relatively small thicknesses. However, they should not be excessively thin to the point of causing significant deformations and are widely used in various aerospace, architectural, civil, marine, and mechanical engineering areas as load-carrying segments [4-6]. The existence of shell structures in these applications induced researchers to understand their behavior to increase efficiency and avoid failure [7]. Therefore, several theories have been published so far on different types of material. Also,

various methods, including analytical, experimental, and numerical techniques, were studied to examine the dynamic response of composite structures [8-10].

Based on first-order shear deformation theory, Javed et al. [11] investigated the dynamics of non-uniform cylindrical shells. An E-glass/epoxy (EGE) composite laminate shell was constructed from Kevlar-49/epoxy composite materials arranged in different orders. Several anti-symmetric angle ply layers were arranged in layers under different boundary conditions to achieve a shell with a non-uniform thickness. Helloty et al. [12] examined the impact of laminated shell composite materials and boundary conditions on the dynamic response of structures under different loads using the finite element method. Chaubey et al. [13] developed a formulation of finite elements that incorporates the third-order shear deformation theory (TSDT) to calculate the non-dimensional natural frequency of composite shells. This formulation also takes into account the presence of concentrated mass and cutouts in the shell structure. The suggested model provides zero transverse shear stress conditions at the shell structure in the head and base without employing a shear correction factor.

Under extreme boundary conditions, Li et al. [14] proposed a new semi-analytical method for solving the influence of free vibration problem on functionally graded porous (FGP) shell composite by utilizing the energy method and first-order shear deformation theory (FSDT). Two typical types of porosity distributions employed in the study were symmetric and nonsymmetrical models. Wang et al. [15] performed free vibration behavior on a metal foam composite cylindrical shell. They specifically focused on the assumed governing variables that dictate the system characteristics i.e., the natural frequency and mode shapes. Dongze He [16] employed a wave-based method to analyze the free vibration of composite laminated cylindrical shells. The study considered shells with both restrained elastic boundaries and general classical boundaries. The analysis aimed to understand the vibration behavior of these shells under various conditions. Ghasemi et al. [17] utilized the modified couple stress theory (MCST) to investigate the frequencies of micro and nano-fiber-metal laminate (FML) cylindrical shells. This study focused on analyzing the vibration characteristics and behavior of FML cylindrical shells using the MCST approach.

Using the finite element method (FEM), Tran et al. [18] investigated laminated graphene nanoplatelets reinforced composite shells. In this paper, using first-order shear deformation theory, an 8-node iso-parametric element has been used to determine the oscillation equation for shell structure. According to Zhang et al. [19], sandwich shells resting on the Pasternak foundation exhibit various vibrational behaviors under humid conditions. A composite laminated spherical shell was examined using Fourier series discretization

techniques [20]. Experimental, analytical, and ABAQUS/CAE analyses were conducted by M Q. Wu et al. [21] to examine the nonlinear dynamic response of two-dimensional square shallow composite laminated shells reinforced by carbon fiber under foundation excitation.

Several laminated composite shells and spatial structures were examined by conducting free vibration analysis by Chen et al. [22]. Utilized the moving-least-squares approximation to develop a mesh-free FSDT method. FG carbon nanotube-reinforced composite panels with sinusoidal corrugation were studied using iso-geometric analysis by Muhammadi et al. [23]. Ameer Melaibari et al. [24] applied a novel analytical technique based on Galerkin's theory in the study of composite laminated shells reinforced with single-walled carbon nanotubes at random orientations. Five types of shell shapes with different geometrical properties are considered in this study. Kuo Tian et al. [25] used a numerical vibration correlation method and optimization design to improve the buckling load of composite cylindrical shells loaded axially. Lin et al. [26] investigated the dynamic stability of cylindrical fiber composite shells with metal liners under uniform pressure pulses by combining finite element simulations and theoretical modeling. Using a semi-analytical approach, Haichao Li [27] studied the free vibrations of laminated composite cylindrical and spherical shells subjected to complex boundary conditions using multi-segment partitioning strategies. According to Sandipan Nath Thakur [28], moderately thick hyperbolic paraboloidal laminated shells exhibit effective response behavior in the presence of higher-order shear deformation theory (HSDT). Donnell's shell theory has been combined with the refined Halpin-Tsai micromechanical approach in a study conducted by Zamani [29]. Laminated composite conical shells reinforced by graphene sheets were investigated for their free vibration behaviors. Using the finite element method (FEM), K. Kim et al. [30] studied elliptical-cylindrical-elliptical laminated composite shells with elastic boundary conditions. A stepped and stiffened cylindrical shell structure subjected to arbitrary boundary conditions was analyzed using stepped and stiffened shear deformation theory and the finite element method [31]. A study by Logesh, et al. [32] evaluated the effect of nanoparticles ($Al_2O_3-SiO_2$) on the mechanical properties of blended matrix polymer composites.

Free vibration of porous graphene nano-platelet (GNP) nano-enrichment composite spheroidal-cylindrical shells was performed by estimating the properties of the composite by using the rule of mixture and Halpin-Tsai homogenization techniques [33]. Free vibration studies on shell structures reinforced with SiO_2 were performed by employing the finite element method and semi-analytical approach [34, 35]. From the host of literature, it was found that the rule of mixture is inadequate in estimating the material properties effectively [36,

37]. Accurate estimation of material properties, such as Young’s modulus and density, is critical as they serve as inputs for analytical and finite element analysis. Incorrect estimation of these properties can lead to erroneous predictions such as natural frequency in the present study. Therefore, one of the key objectives of the present work is to experimentally determine the material properties and use these values to predict the natural frequency of Al₂O₃ and SiO₂-reinforced structures using established methods.

Additionally, the shell structure can be used as an orthosis or prosthetic part; thus, it was necessary to incorporate reinforcement into the shell structure to determine its mechanical characteristics. Nanoparticles can be used as reinforcement materials for enhancing the mechanical and dynamic characteristics. Then, the reinforcement of nanoparticle materials changes the stiffness-to-weight ratio. So, the new contribution of this work is to modify the mechanical and dynamic characterizations of shell structure by reinforcing it with nanoparticle materials to achieve high strength and high mechanical and dynamic performance. Various interesting findings are presented here concerning the static and dynamic analysis of nanocomposite laminated thin shell structure. A unified formulation to investigate the vibration behaviors of cylindrical shells subjected to arbitrary boundary restraints was performed. This formulation is based on various parameters such as volume fraction, nanoparticle type, and boundary conditions.

2. MATERIALS AND METHODS

In the literature, numerous researchers employed different experimental techniques focusing on the behavior of nanocomposite structures [38, 39]. In the current study, the experiment program consists of two parts: the first one is preparing samples with and without nanomaterials (Al₂O₃ & SiO₂), and the second part is related to obtaining Young’s modulus by performing tensile tests. The samples were produced from the fibers materials (Perlon, Carbon, Kevlar, Kenaf, Resin, Hardener, and two types of nanoparticles) as illustrated in Figure 1. The specimens were created in the form of layering reinforced materials that were pressed for each sample in groups A and B, as shown in Table 1-3. Figure 1 shows photographic images of the material used and the packaging of Orthocryl Resin and powders of nanoparticles, displaying the product name and brand. This allows researchers and industries to easily procure the product and access important details, such as purity, density, moisture content, chemical composition, safety data, and usage guidelines, among other key specifications.

All the produced specimens are reinforced with two types of nanoparticles with varied volume fractions. Figure 2 illustrates how the production of

the combined model with bonding nanomaterials influences the mixing of nanoparticles with thickening liquid in orthocryl resin using an ultrasonic homogenizer device. Figure 3 shows how reinforcement is accomplished by mixing resin and layers of fibers using the vacuum technique. Nanoparticles utilized to improve composite materials are SiO₂ and Al₂O₃ with five values of volume fractions (0.5, 1, 1.5, 2, and 2.5) %.

Young’s modulus of elasticity of composite materials with different nanomaterials was determined using the Universal Tensile Machine by



Fig. 1. Materials used in preparing the sample



Fig. 2. Vacuum machine

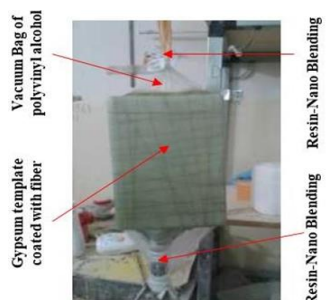


Fig. 3. Materials Manufacturing

following the standard specification (ASTM-D638), as shown in Figure 4. Table 2 shows the ten samples with each volume fraction that were tested for tensile strength.

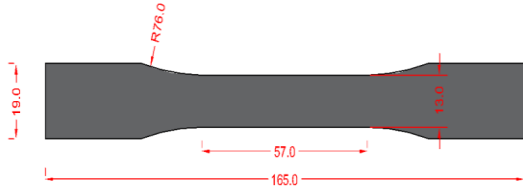


Fig. 4. Geometry of tensile test Specimen According to (ASTM-D638) standard

Table 1. Mechanical properties of composite material used

Group	Number of Layers	E (GPa)	ν	ρ (kg/m ³)
A	1 layer Perlon	16.78	0.22	1180
	2 layer Carbon			
	1 layer Kevlar			
	4 layer Perlon			
	1 layer Kevlar			
	2 layer Carbon			
B	1 layer Perlon	17.49	0.18	1210
	2 layer Carbon			
	1 layer Kevlar			
	1 layer Kenaf			
	2 layer Perlon			
	1 layer Kenaf			
B	1 layer Kevlar	17.49	0.18	1210
	2 layer Carbon			
	1 layer Perlon			
	1 layer Perlon			
	2 layer Carbon			
	1 layer Perlon			

Table 2. Material properties used with different volume fractions of nanomaterials

Group	Nano Al ₂ O ₃ Material			Nano SiO ₂ Material		
	S	V _f %	E (GPa)	S	V _f %	E (GPa)
A	A1	0	16.78	A1	0	16.78
	A2	0.5	17.34	A2	0.5	17.92
	A3	1	18.42	A3	1	19.12
	A4	1.5	19.83	A4	1.5	20.33
	A5	2	20.88	A5	2	21.74
	A6	2.5	22.47	A6	2.5	23.87
B	B1	0	17.49	B1	0	17.49
	B2	0.5	18.23	B2	0.5	19.46
	B3	1	19.47	B3	1	20.74
	B4	1.5	20.44	B4	1.5	22.31
	B5	2	21.89	B5	2	23.85
	B6	2.5	24.13	B6	2.5	25.78

Table 3. Details of composite constituents

Item	Group A	Group B
layer Perlon	6	4
layer Carbon	4	4
layer Kevlar	2	2
layer Kenaf	0	2

3. THE ANALYTICAL MODEL

The analytical solutions of the shell structure include constructing an equation of motion for a

simply supported shell structure. However, to derive the general vibration shell equation, the following steps are required:

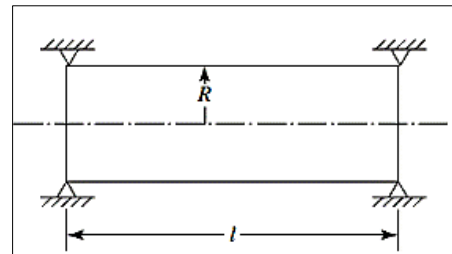
- Derive the equation of motion for the shell structure considering external load effects.
- Determine the generalized displacement for the shell by assuming it is a function of the selected coordinates.
- Calculate the natural frequency using the derived mathematical model.

In this context, the cartesian system coordinates are utilized [34].

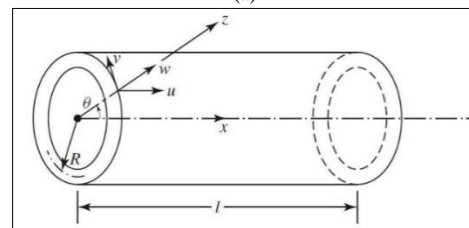
$$\begin{aligned} X &= X(\alpha, \beta) \\ Y &= Y(\alpha, \beta) \\ Z &= Z(\alpha, \beta) \end{aligned} \quad (1)$$

The two independent coordinates α and β can be used to describe the undeformed middle surface of a thin shell; therefore, the following assumptions for small displacement theory apply as follows:

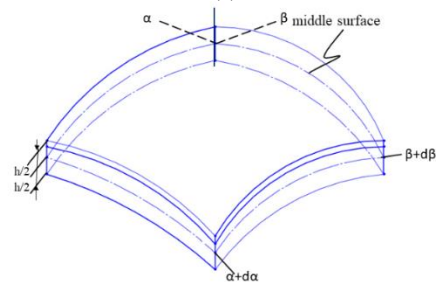
- Small thickness compared to other shell dimensions.
- Minimal shell strain and displacement.
- Negligible stress in the z-direction.
- The normal to the undeformed middle surface does not undergo any extension or contraction, even after deformation. It maintains its straightness and remains perpendicular to the middle surface throughout the deformation process. Then, the displacement representation of the shell is assumed as [40,41]:



(a)



(b)



(c)

Fig. 5. The model's geometry used (a) Simply-supported BCs and (b-c) coordinates of a thin cylindrical shell

$$\begin{aligned} \bar{u}(\alpha, \beta, z, t) &= u(\alpha, \beta, t) + z\theta_\alpha(\alpha, \beta, t) \\ \bar{v}(\alpha, \beta, z, t) &= v(\alpha, \beta, t) + z\theta_\beta(\alpha, \beta, t) \\ \bar{w}(\alpha, \beta, z, t) &= w(\alpha, \beta, t) \end{aligned} \quad (2)$$

The variables (u, v, and w) represent the shell's middle surface displacements along the directions α , β , and z, while R indicates the radius of the cylinder. The rotations of the intermediate surface normal about the α and β axes, denoted as θ_α and θ_β , can be determined as follows:

$$\begin{aligned} \theta_\alpha &= \frac{\partial w}{\partial x} \\ \theta_\beta &= \frac{v}{R} - \frac{1}{R} \frac{\partial w}{\partial \theta} \end{aligned} \quad (3)$$

For the cylindrical shell shown in Figure 5, with coordinates x, θ , and z, the parallel displacements along the cylindrical shell coordinate are denoted as u, v, and w, while the variables α and β represent x and θ , respectively, and h is the thickness of the shell. Thus, the general equation of motion for a cylindrical shell can be expressed as follows:

$$\begin{aligned} \left(\frac{\partial N_{xx}}{\partial x} + \frac{1}{R} \frac{\partial N_{\theta x}}{\partial \theta} + f_x \right) &= \rho h \ddot{u} \\ \frac{\partial N_{x\theta}}{\partial x} + \frac{1}{R} \frac{\partial N_{\theta\theta}}{\partial \theta} + \frac{1}{R} \left(\frac{\partial M_{x\theta}}{\partial x} + \frac{1}{R} \frac{\partial M_{\theta\theta}}{\partial \theta} \right) + f_\theta &= \rho h \ddot{v} \\ \frac{\partial^2 M_{xx}}{\partial x^2} + \frac{1}{R} \frac{\partial^2 M_{x\theta}}{\partial x \partial \theta} + \frac{1}{R^2} \frac{\partial^2 M_{\theta\theta}}{\partial \theta^2} - \frac{N_{\theta\theta}}{R} + f_z &= \rho h \ddot{w} \end{aligned} \quad (4)$$

where, N_{xx} , $N_{\theta x}$, $N_{\theta\theta}$, M_{xx} , $M_{x\theta}$, and $M_{\theta\theta}$ are force and moment components, shown in Figure 6. The components are written as:

$$\begin{aligned} N_{xx} &= C \left(\frac{\partial u}{\partial x} + \frac{v}{R} \frac{\partial v}{\partial \theta} + \frac{v}{R} \right) \\ N_{\theta x} &= C \frac{(1-\nu)}{2} \left(\frac{\partial v}{\partial x} + \frac{1}{R} \frac{\partial u}{\partial \theta} \right) \\ N_{\theta\theta} &= C \left(\frac{1}{R} \frac{\partial v}{\partial \theta} + \frac{w}{R} + v \frac{\partial u}{\partial x} \right) \\ M_{xx} &= D \left(-\frac{\partial^2 w}{\partial x^2} + \frac{v}{R^2} \frac{\partial v}{\partial \theta} - \frac{v}{R^2} \frac{\partial^2 w}{\partial \theta^2} \right) \\ M_{x\theta} &= D \frac{(1-\nu)}{2} \left(\frac{1}{R} \frac{\partial v}{\partial x} - \frac{2}{R} \frac{\partial^2 w}{\partial x \partial \theta} \right) \\ M_{\theta\theta} &= D \left(\frac{1}{R^2} \frac{\partial v}{\partial \theta} - \frac{1}{R^2} \frac{\partial^2 w}{\partial \theta^2} - v \frac{\partial^2 w}{\partial x^2} \right) \end{aligned} \quad (5)$$

where, $C = \frac{Eh}{(1-\nu^2)}$, and $D = \frac{Eh^3}{12(1-\nu^2)}$.

Shell materials have a modulus of elasticity (E) corresponding to their thickness and density, respectively; they also have a Poisson's ratio (ν) corresponding to the material, while external forces are (f_x, f_y, f_z).

Then, by inserting Equation (5) into Equation (4), for free vibration external forces $f_x, f_y, f_z = 0$, and by using Donnell-Mushtari Vlasov (DMV) theory, get

$$\begin{aligned} \frac{\partial^2 u}{\partial x^2} + \frac{(1-\nu)}{2R^2} \frac{\partial^2 u}{\partial \theta^2} + \frac{v}{R} \frac{\partial w}{\partial x} + \frac{(1+\nu)}{2R} \frac{\partial^2 v}{\partial x \partial \theta} &= \frac{(1-\nu^2)}{E} \rho \frac{\partial^2 u}{\partial t^2} \\ \left(\frac{(1-\nu)}{2} \frac{\partial^2 v}{\partial x^2} + \frac{1}{R^2} \frac{\partial^2 v}{\partial \theta^2} + \frac{1}{R^2} \frac{\partial w}{\partial \theta} + \frac{(1+\nu)}{2R} \frac{\partial^2 u}{\partial x \partial \theta} \right) &= \frac{(1-\nu^2)}{E} \rho \frac{\partial^2 v}{\partial t^2} \\ \left(-\left(\frac{v}{R} \frac{\partial u}{\partial x} + \frac{1}{R^2} \frac{\partial v}{\partial \theta} + \frac{w}{R^2} \right) - \frac{h^2}{12} \left(\frac{\partial^4 w}{\partial x^4} + \frac{2}{R^2} \frac{\partial^4 w}{\partial x^2 \partial \theta^2} + \frac{1}{R^4} \frac{\partial^4 w}{\partial \theta^4} \right) \right) &= \frac{(1-\nu^2)}{E} \rho \frac{\partial^2 w}{\partial t^2} \end{aligned} \quad (6)$$

In Equation 6, (F) represents the bending load, (l) denotes the span length, (A) is the area of the beam's cross-section, and (G) represents the modulus of shear.

Therefore, with the aid of Equation 6 and the separation of variable technique, the shell

displacement as a function of x and θ can be found. The following relations can be used for simply supported boundary conditions ($x = 0$ and $x = l$).

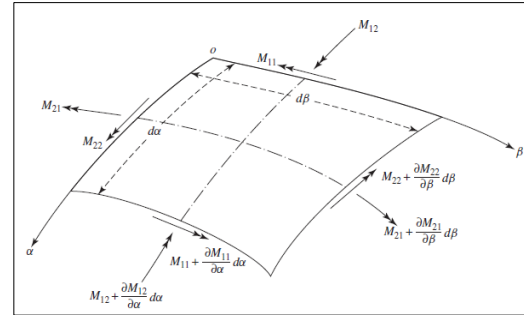
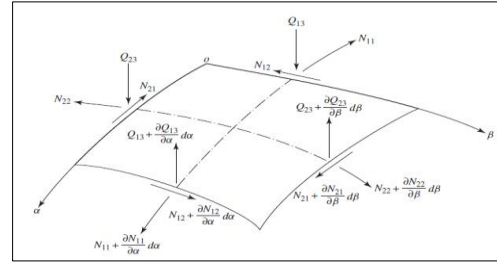


Fig. 6. (a) A shell force resultant (b) A resultant of the moment

$$\begin{aligned} v(0, \theta, t) &= w(0, \theta, t) = 0 \\ N_{xx}(0, \theta, t) &= M_{xx}(0, \theta, t) = 0 \\ v(l, \theta, t) &= w(l, \theta, t) = 0 \\ N_{xx}(l, \theta, t) &= M_{xx}(l, \theta, t) = 0 \end{aligned} \quad (7)$$

Then, the components of displacement of the cylindrical shell can be represented by assuming the general form of harmonic function:

$$\begin{aligned} u(x, \theta, t) &= \sum_m^\infty \sum_n^\infty A_{mn} \times \cos \frac{m\pi x}{l} \times \cos n\theta \times \cos \omega t \\ v(x, \theta, t) &= \sum_m^\infty \sum_n^\infty B_{mn} \times \sin \frac{m\pi x}{l} \times \sin n\theta \times \cos \omega t \\ w(x, \theta, t) &= \sum_m^\infty \sum_n^\infty C_{mn} \times \sin \frac{m\pi x}{l} \times \cos n\theta \times \cos \omega t \end{aligned} \quad (8)$$

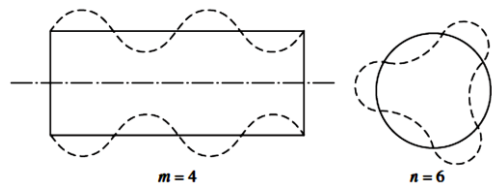


Fig.7. Displacement mode of a Shell

As depicted in Figure 7, m and n denote the wave characteristics of displacement along the length and circumference of the shell, respectively. By substituting Equation (8) into Equation (6), the following form was obtained:

$$\begin{aligned} &\left(-\left(\frac{m\pi}{l} \right)^2 A_{mn} \cos \frac{m\pi x}{l} \times \cos n\theta \right) - \\ &\left(\frac{(1-\nu)}{2R^2} (n)^2 A_{mn} \cos \frac{m\pi x}{l} \times \cos n\theta \right) + \\ &\left(\frac{v}{R} \left(\frac{m\pi}{l} \right) C_{mn} \cos \frac{m\pi x}{l} \times \cos n\theta \right) \\ &+ \left(\frac{(1+\nu)}{2R} \left(\frac{m\pi}{l} \right) (n) B_{mn} \cos \frac{m\pi x}{l} \times \cos n\theta \right) \\ &\left(-\frac{(1-\nu^2)}{E} \rho \omega^2 A_{mn} \cos \frac{m\pi x}{l} \times \cos n\theta \right) \cos \omega t = \end{aligned}$$

$$\begin{pmatrix}
 -\left(\frac{1-\nu}{2}\right)\left(\frac{m\pi}{l}\right)^2 B_{mn} \times \sin \frac{m\pi x}{l} \times \sin n\theta \\
 -\left(\frac{1}{R^2}\right)(n)^2 B_{mn} \times \sin \frac{m\pi x}{l} \times \sin n\theta - \\
 \left(\frac{1}{R^2}\right)(n) C_{mn} \times \sin \frac{m\pi x}{l} \times \sin n\theta \\
 +\left(\frac{1+\nu}{2R}\right)\left(\frac{m\pi}{l}\right)(n) A_{mn} \times \sin \frac{m\pi x}{l} \times \sin n\theta \\
 \left(-\frac{(1-\nu^2)}{E}\rho\omega^2 B_{mn} \times \sin \frac{m\pi x}{l} \times \sin n\theta\right) \cos \omega t \\
 \left(\frac{\nu}{R}\left(\frac{m\pi}{l}\right) A_{mn} \sin \frac{m\pi x}{l} \times \cos n\theta\right) - \\
 \left(\frac{1}{R^2}\right)(n) B_{mn} \sin \frac{m\pi x}{l} \times \cos n\theta - \\
 \left(\frac{1}{R^2}\right) C_{mn} \sin \frac{m\pi x}{l} \times \cos n\theta - \\
 \left(\frac{h^2}{12}\right)\left(\frac{m\pi}{l}\right)^4 C_{mn} \sin \frac{m\pi x}{l} \times \cos n\theta - \\
 \left(\frac{h^2}{12R^2}\right)\left(\frac{m\pi}{l}\right)^2 (n)^2 C_{mn} \times \sin \frac{m\pi x}{l} \times \cos n\theta \\
 -\left(\frac{h^2}{12R^4}\right)(n)^4 C_{mn} \times \sin \frac{m\pi x}{l} \times \cos n\theta \\
 \left(-\frac{(1-\nu^2)}{E}\rho\omega^2 C_{mn} \times \sin \frac{m\pi x}{l} \times \cos n\theta\right) \cos \omega t
 \end{pmatrix} \cos \omega t = \begin{pmatrix}
 \left(\frac{m\pi}{l}\right)^2 + \frac{(1-\nu)}{2R^2}(n)^2 \\
 -\frac{(1-\nu^2)}{E}\rho\omega^2 \\
 \left(-\frac{(1+\nu)}{2R}\left(\frac{m\pi}{l}\right)(n)\right) \\
 \left(\frac{1-\nu}{2}\left(\frac{m\pi}{l}\right)^2 + \frac{1}{R^2}(n)^2\right) \\
 -\frac{(1-\nu^2)}{E}\rho\omega^2 \\
 \left(-\frac{\nu}{R}\left(\frac{m\pi}{l}\right)\right) \\
 \left(\frac{1}{R^2}(n)\right) \\
 \left(\frac{1}{R^2} + \frac{h^2}{12}\left(\frac{m\pi}{l}\right)^4 + \frac{h^2}{12R^2}\left(\frac{m\pi}{l}\right)^2 (n)^2 + \frac{h^2}{12R^4}(n)^4 - \frac{(1-\nu^2)}{E}\rho\omega^2\right) \\
 \left(\frac{1}{R^2}(n)\right) \\
 \left(\frac{1}{R^2}(n)\right)
 \end{pmatrix} \begin{pmatrix} A_{mn} \\ B_{mn} \\ C_{mn} \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad (9)$$

$$\begin{pmatrix}
 \left(\frac{m\pi}{l}\right)^2 + \frac{(1-\nu)}{2R^2}(n)^2 \\
 -\frac{(1-\nu^2)}{E}\rho\omega^2 \\
 \left(-\frac{(1+\nu)}{2R}\left(\frac{m\pi}{l}\right)(n)\right) \\
 \left(\frac{1-\nu}{2}\left(\frac{m\pi}{l}\right)^2 + \frac{1}{R^2}(n)^2\right) \\
 -\frac{(1-\nu^2)}{E}\rho\omega^2 \\
 \left(-\frac{\nu}{R}\left(\frac{m\pi}{l}\right)\right) \\
 \left(\frac{1}{R^2}(n)\right) \\
 \left(\frac{1}{R^2} + \frac{h^2}{12}\left(\frac{m\pi}{l}\right)^4 + \frac{h^2}{12R^2}\left(\frac{m\pi}{l}\right)^2 (n)^2 + \frac{h^2}{12R^4}(n)^4 - \frac{(1-\nu^2)}{E}\rho\omega^2\right) \\
 \left(\frac{1}{R^2}(n)\right) \\
 \left(\frac{1}{R^2}(n)\right)
 \end{pmatrix} \begin{pmatrix} A_{mn} \\ B_{mn} \\ C_{mn} \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad (10)$$

Or, Equation 11 may be arranged as,

$$\begin{pmatrix}
 \left(\frac{m\pi}{l}\right)^2 + \frac{(1-\nu)}{2R^2}(n)^2 \\
 -\frac{(1-\nu^2)}{E}\rho\omega^2 \\
 \left(-\frac{(1+\nu)}{2R}\left(\frac{m\pi}{l}\right)(n)\right) \\
 \left(\frac{1-\nu}{2}\left(\frac{m\pi}{l}\right)^2 + \frac{1}{R^2}(n)^2\right) \\
 -\frac{(1-\nu^2)}{E}\rho\omega^2 \\
 \left(-\frac{\nu}{R}\left(\frac{m\pi}{l}\right)\right) \\
 \left(\frac{1}{R^2}(n)\right) \\
 \left(\frac{1}{R^2} + \frac{h^2}{12}\left(\frac{m\pi}{l}\right)^4 + \frac{h^2}{12R^2}\left(\frac{m\pi}{l}\right)^2 (n)^2 + \frac{h^2}{12R^4}(n)^4 - \frac{(1-\nu^2)}{E}\rho\omega^2\right) \\
 \left(\frac{1}{R^2}(n)\right) \\
 \left(\frac{1}{R^2}(n)\right)
 \end{pmatrix} \begin{pmatrix} A_{mn} \\ B_{mn} \\ C_{mn} \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad (11)$$

Then, by solution, Equation 12 may be used to estimate the natural frequency for a model α , β , and z -direction, with various modulus of

elasticity value and different nanoparticle materials effect.

4. NUMERICAL SIMULATION

With Finite Element Methods (FEMs), numerical findings can be compared with those gained through experiments. The ANSYS software 2021 R1 was used to perform a modal analysis of a composite cylindrical structure. Using the Solid186 shell model, 1692 shell elements were generated, and meshing was performed, as shown in Figure 8; shell end conditions were selected, and modal analyses were conducted for each desired model [42-45]. The experimental work yields the mechanical properties used in the engineering data view. Analytical techniques (Equation 12) and numerical approaches (FEM) can be used to calculate the natural frequency of shell structures. Table 1 shows two groups composed of different layers of composite materials reinforced with nanoparticles.

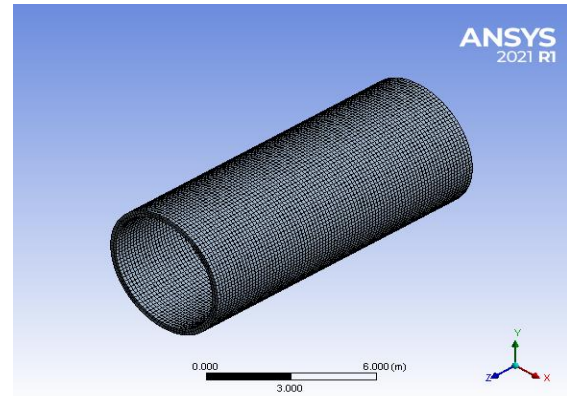


Fig. 8. A 3D model with a mesh

5. RESULTS AND DISCUSSION

The main objective of this investigation is to modify the mechanical vibration characteristics of composite shell structures with a high volume fraction of the nanoparticle materials. To enhance the mechanical performance of such structures to be used in engineering applications. The influence of reinforced nanoparticles in the composite shell (Nano Al_2O_3 and Nano SiO_2) is determined using analytical and numerical approaches.

Adding two layers of Kenaf to group B at the expense of Perlon layers in group A led to an increase in Young's modulus, meaning good material strength of composite samples from (16.78 to 17.49) GPa with a percentage increase of 4.23%. The addition of the volume fraction of nano Al_2O_3 and nano SiO_2 in groups A and B of the composite shell at value (0-2.5)% gave a direct relationship. The higher the volume fraction, the better the resistance of the composite material, the improved the modulus of elasticity (22.47 and 24.13) GPa of Al_2O_3 and (23.87 and 25.78) GPa of SiO_2 at 2.5% volume fraction.

Table 4 shows that adding SiO₂ nanoparticles had the highest modulus of elasticity (23.87 and 25.78) GPa of SiO₂ compared to (22.47 and 24.13) GPa of Al₂O₃ nanoparticles. As a result, as shown in Table 4, the inclusion of SiO₂ nanoparticles in group B enhances the modulus of elasticity by 47.39 percent at a volume fraction of 2.5 percent.

Table 4. Variation in modulus of elasticity with and without reinforcements

Group	E (GPa)	E (GPa)	Incr.	E	Incr.
	Without Nano	Al ₂ O ₃ Nano	%	(GPa) SiO ₂ Nano	%
A	16.78	22.47	33.90	23.87	42.25
B	17.49	24.13	37.96	25.78	47.39

Figure 9 demonstrates the effect of adding nanoparticles SiO₂ and Al₂O₃ to the composite shell on the elastic modulus. Based on the findings, groups A and B show good agreement between the numerical and analytical natural frequencies. Furthermore, it has been observed that the presence of SiO₂ nanoparticles increases the natural frequency to a greater extent than the effect of Al₂O₃ nanoparticles in both groups. The possible reason is due to the nature of the microstructure of Si and Al. However, group B recorded the highest natural frequency of 5780 rad/sec for SiO₂ nano particles at a volume fraction of 2.5%. Figure 10 shows the effect of nanoparticle reinforcements on composite shell structures by comparing analytical and numerical frequency results.

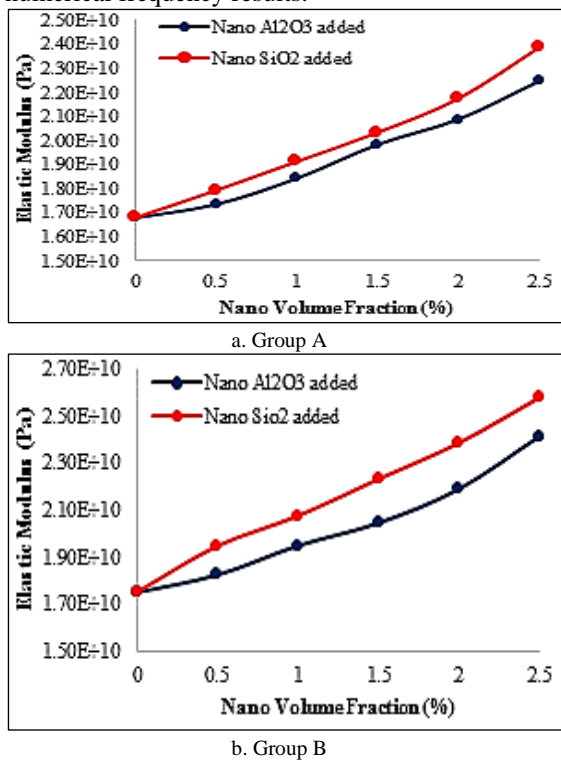
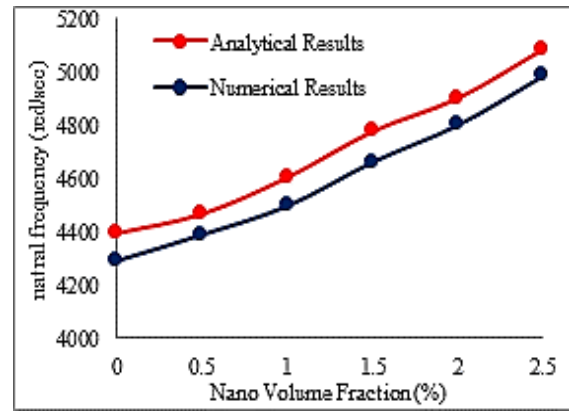
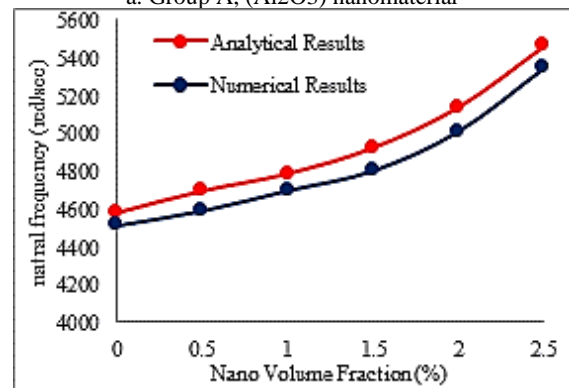


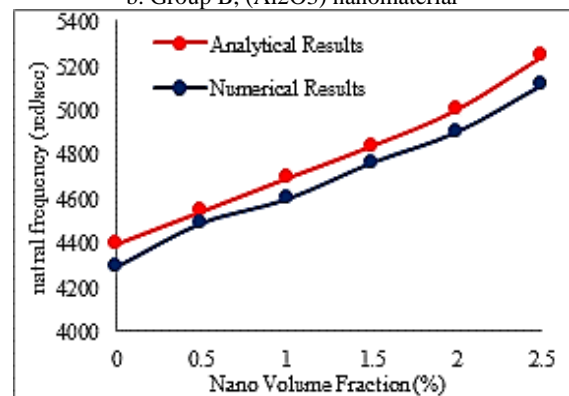
Fig. 9. Various composite shell groups exhibit different elastic modulus variations with nano volume fractions



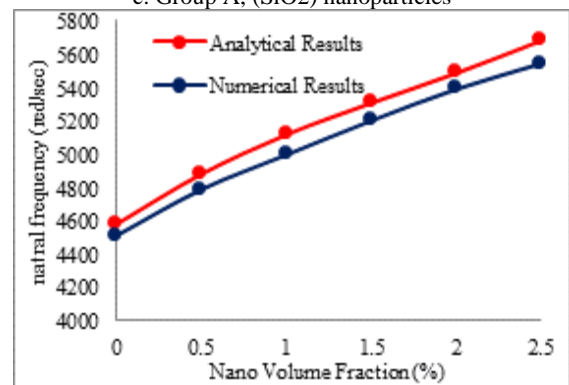
a. Group A, (Al₂O₃) nanomaterial



b. Group B, (Al₂O₃) nanomaterial



c. Group A, (SiO₂) nanoparticles

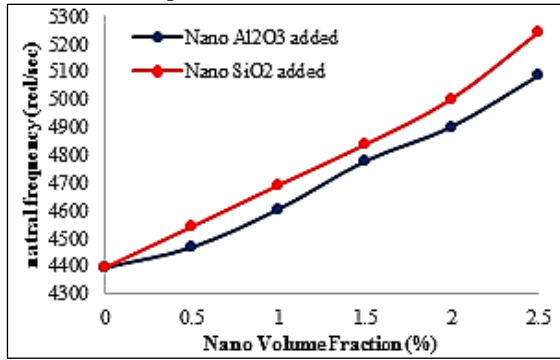


d. Group B, (SiO₂) nanoparticles

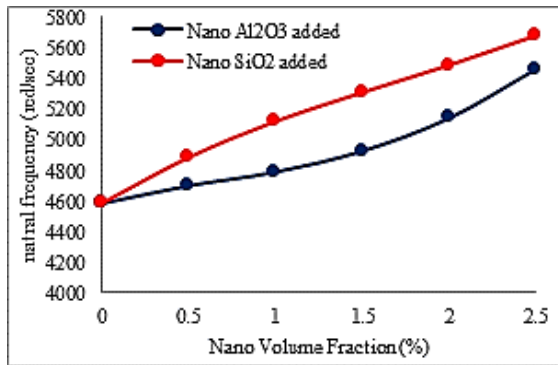
Fig. 10. The impacts of nanoparticle materials on composite shell structures by comparing analytical and numerical frequency results

Figure 11 shows the numerical and analytical results of the natural frequency for two groups (A and B) based on the concentration of both

nanoparticles, SiO₂ and Al₂O₃. The results show a clear improvement in the natural frequency of shells with SiO₂ nanoparticles rather than Al₂O₃.



a. Group A



b. Group B

Fig. 11. Free vibration analysis of the composite shell types A and B with different nano volume fractions

Normalized natural frequency results have been obtained when studying the free vibration of a cylindrical shell with varied BCs [30]:

$$\psi = \omega \cdot R \sqrt{\rho(1 - \nu^2)/E} \tag{13}$$

Figure 12 illustrates the numerical outcomes of the frequency parameter for the composite shell group (A) at a ratio of (L/R = 5) with five different boundary conditions (BCs). Based on the obtained results, it is evident that the frequency coefficient value rises as the number of restrictions in the chosen model increases. This finding indicates an increase in the overall strength of the shell model.

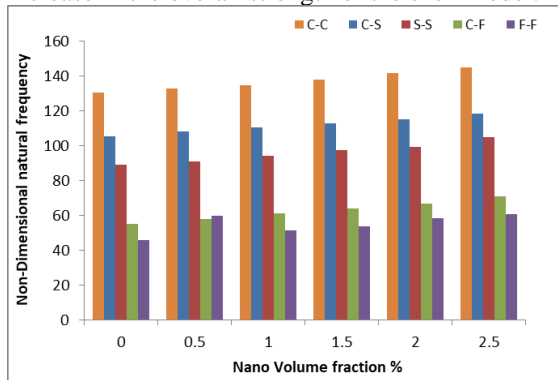
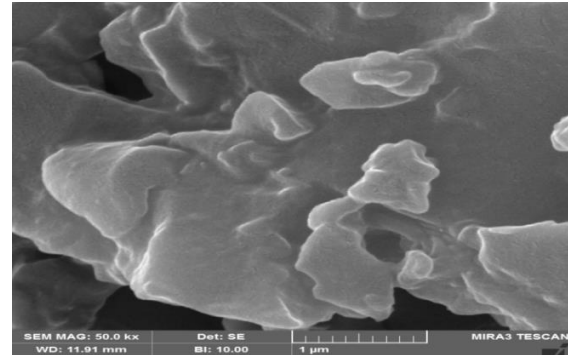


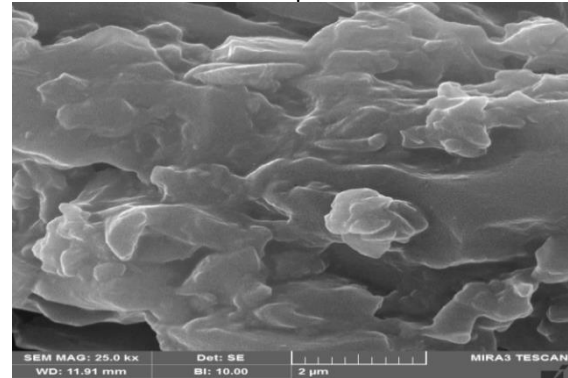
Fig. 12. Results for the frequency parameter (Group A) subjected to various BCs

SEM has been widely used to examine surface topology and nanoparticles. Different materials containing nanoparticles of known distribution were analyzed to test the accuracy of an analytical model using high-resolution SEM. Figure 13 shows SEM images of Group A with Nano SiO₂ with different magnifications (1 μm and 2 μm). Figures 14-16 show SEM images of Group A with nano Al₂O₃ and SiO₂ material with different magnifications, respectively. From the results, one can see the distribution of fibers into the matrix phase. Different magnifications of SEM images are presented that illustrate the reinforcement distribution and binding of phases.

Moreover, it can be seen that SiO₂ images exhibit stronger borders and more contrast than Al₂O₃. Also, it is observed that a good transmission contrast is observed between the two types of nanoparticles. High-resolution SEM imaging can define nanoparticles' morphology, characteristics, and interior structure.

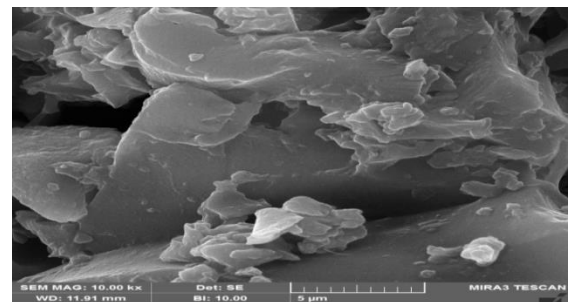


a. 1 μm

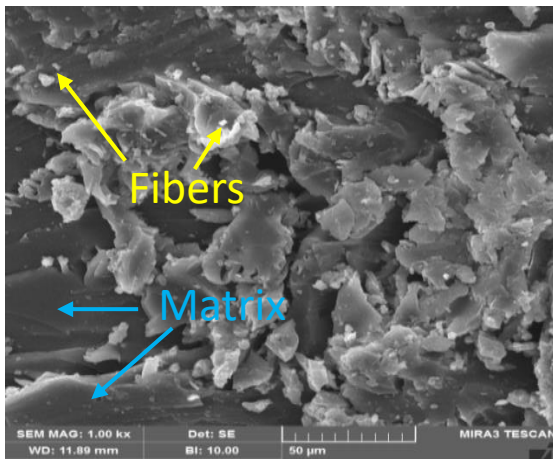
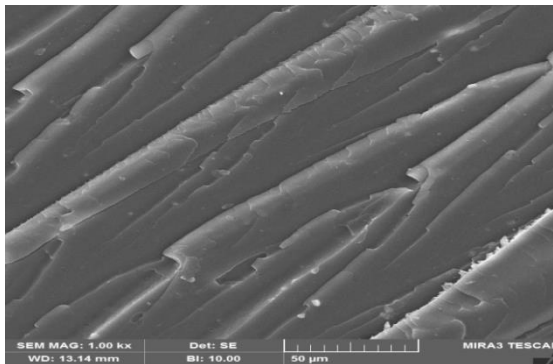
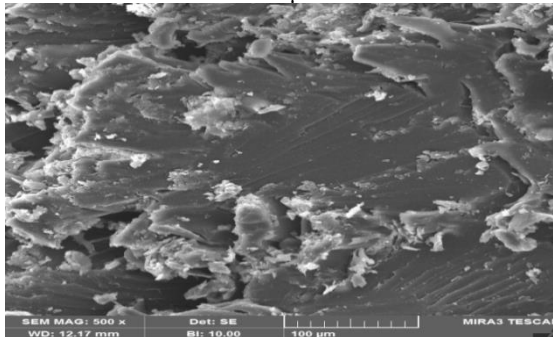
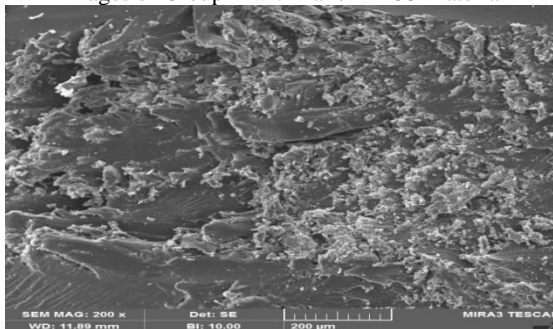
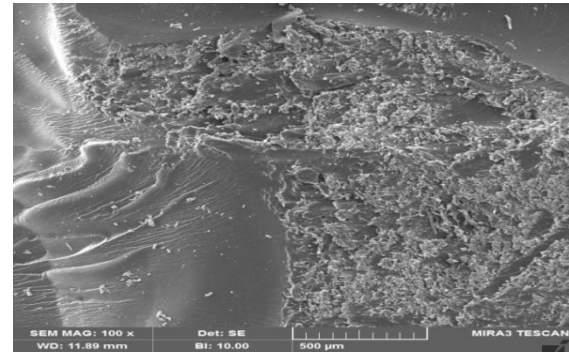


b. 2 μm

Fig. 13. SEM images of Group A with Nano SiO₂ with different magnifications



a. 5 μm

b. 50 μm Fig. 14. SEM images of Group A with nano Al_2O_3 Material with different magnificationa. 50 μm b. 100 μm Fig. 15. Different magnifications of SEM images of Group B with nano Al_2O_3 Materialb. 200 μm b. 500 μm Fig. 16. Different magnifications of SEM images of Group B with nano SiO_2 Material

4. CONCLUSION

The addition of nano Al_2O_3 and Nano SiO_2 reinforced composite cylindrical shells for groups A and B is covered in detail in this work. There are several interpretations can be recorded to sustain the nanocomposite shell, mechanical and natural frequency properties, including:

1. The modulus elasticity values of the composite shell material in groups A and B were increased by adding small nano-volume fractions of nanoscale Al_2O_3 and nanoscale SiO_2 nanoparticles.
2. The molecular bonds between the shell composite material particles are strengthened by the nanoparticles that permeate between them, contributing to the increased strength of the two groups' composite shell material.
3. The addition of nano Al_2O_3 and nano SiO_2 nanoparticles increases the natural frequency of the shell due to the sitting of nanoparticles through voids of composite materials. This increases the bonding strength between particles and grains in the atomic structure of the composite material's shell. Two composite materials were strengthened by the addition of nanoparticle materials, resulting in a roughly 30% increase in natural frequency.
4. Adding nano volume fraction to the composite materials in a shell increases the modulus of elasticity values which in turn increases the natural frequency of a composite.
5. By comparing the results, adding nano SiO_2 increases the modulus of elasticity, and vibration characteristics that increase the shell's strength better than nano Al_2O_3 .
6. Free vibration analysis by employing the analytical and numerical techniques shows good agreement between each other with a maximum difference of 5%.

Source of funding: *This research received no external funding.*

Author contributions: *Research concept and design, E.K.; Collection and/or assembly of data, Z.A.;*

Writing the article, M.A.; Critical revision of the article, R.M.; Final approval of the article, E.N.

Declaration of competing interest: *The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.*

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